Emergence of geometry in stochastic systems

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motivation ³analogies between quantum systems and stochastic systems [Nelson(1966), Parisi,Wu(1981)] Can quantum gravity be formulated by using stochastic processes?

main results [We here focus on Markov chains as stochastic processes.]

• For a given Markov chain, we introduce a geometry in configuration space by defining the distance between configurations.

• We show that [AdS geometry emerges for the simulated tempering algorithm. [Marinari, Parisi(1992)] we can optimize parameters in numerical simulations geometrically. (- practical use of the distance)

(A) Definition of the distance

- We first prepare some quantities to define the distance:
 - $\mathcal{M} \equiv \{x\}$: configuration space, S(x): action
 - Markov chain with a transition matrix $P(x|y) \ (\equiv \langle x|\hat{P}|y\rangle)$:

• Examples • <u>quadratic action</u>: $S(x) = \frac{\omega}{2}x^2$ analytic form of the distance: $d_n(x, y) = \sqrt{\frac{\omega}{2\sinh(\omega t)}} |x - y|$ (using the Langevin algorithm) • <u>cosine action</u>: $S(\mathbf{x}; \beta_0) = \beta_0 \sum_{\mu=1}^{D} (1 - \cos(2\pi x_{\mu})) - - - - (\star)$

 $x_0 \xrightarrow{\hat{P}} x_1 \xrightarrow{\hat{P}} x_2 \xrightarrow{\hat{P}} \dots$

Probability distribution after *n* steps: $P_n(x|x_0) \equiv \langle x|\hat{P}^n|x_0\rangle$

-Assumptions on \hat{P} :

- $\int (1) \text{ unique convergence: } P_n(x|x_0) \xrightarrow{n \to \infty} \frac{e^{-S(x)}}{Z} (\forall x_0) \left[Z \equiv \int dx \ e^{-S(x)} \right]$
- (2) detailed balance: $P(x|y)e^{-S(y)} = P(y|x)e^{-S(x)}$

③ eigenvalues are all positive

• Transfer matrix: $\hat{T} \equiv e^{S(\hat{x})/2} \hat{P} e^{-S(\hat{x})/2}$ \hat{T} is positive and symmetric (due to 3 and 2 above).

Spectral decomposition of \hat{T} : $\hat{T} = |0\rangle\langle 0| + \sum_{k\geq 1} \lambda_k |k\rangle\langle k| \quad \left[\langle x|0\rangle = \sqrt{\frac{e^{-S(x)}}{Z}}\right] \left[1 > \lambda_1 \geq \lambda_2 \geq \cdots\right]$

Relaxation to the equilibrium can be understood as the decoupling of higher modes: $\hat{T}^n \xrightarrow{n \to \infty} |0\rangle \langle 0|$.

• Connectivity $f_n(x, y)$:

probability to find a path connecting x and y $f_n(x,y) \equiv |$ in the set of *n*-step paths in equilibrium

n-step paths

 $|x, n/2\rangle \ \theta_n(x, y)$

 $\langle y, n/2 \rangle$



 To investigate the large scale geometry of the configuration space, we coarse-grain the configuration space. (identify configurations in the same mode as a single configuration) e.g. for the cosine action (*)

 $d_n(x,y)$ does satisfy the triangle inequality after coarse-graining.

Metric on
$$\overline{\mathcal{M}}$$
:
 $ds^2 \equiv d_n^2(x, x + dx) = \sum_{\mu, \nu} g_{\mu\nu}(x) dx_{\mu} dx_{\nu}$. $(x, x + dx : \text{nearby points in } \overline{\mathcal{M}})$

(B) Emergence of AdS geometry

•We consider a system with highly multimodal equilibrium distribution and with highly degenerate vacua. (typical action: the cosine action (\star))

Normalized connectivity $F_n(x, y)$:

 $= P_n(x|y) \cdot \frac{e^{-S(y)}}{z}$

$$F_n(x,y) \equiv \frac{f_n(x,y)}{\sqrt{f_n(x,x)f_n(y,y)}} = \frac{\langle x|\hat{T}^n|y\rangle}{\sqrt{\langle x|\hat{T}^n|x\rangle\langle y|\hat{T}^n|y\rangle}}$$

Notes

- $\circ F_n(x,x) = 1$
- With normalized states $|x, n/2\rangle \equiv \frac{\hat{T}^{n/2} |x\rangle}{|||\hat{T}^{n/2} |x\rangle||}$, $F_n(x,y) = \langle x, n/2 | y, n/2 \rangle_{\bullet}$

• We define the distance by

 $\theta_n(x,y) \equiv \arccos F_n(x,y)$

 $\theta_n(x,y)$ is actually an inner angle between the normalized states $|x, n/2\rangle$, $|y, n/2\rangle$.

Basic properties

• $\theta_n(x,y)$ satisfies the axioms of distance.

- •We implement the <u>simulated tempering algorithm</u>: (an algorithm to accelerate relaxation [Marinari, Parisi(1992)])
- We extend the configuration space

from \mathcal{M} to $\mathcal{M} \times \mathcal{A} \equiv \{x, \beta_a\}$. β_a : the overall coefficient of the action β_0 (original value) > $\beta_1 > \cdots > \beta_A$

different modes are now connected in the small β_a region

- extended config sp original config sp $\mathcal{M} imes \mathcal{A}$ transition path /
- To investigate the large scale geometry, we introduce extended, coarse-grained configuration space $\overline{\mathcal{M}} \times \mathcal{A}$.
- One can show that the metric on $\bar{\mathcal{M}} \times \mathcal{A}$ is AdS : $(\lambda R^2$

$$ds^{2} = l^{2} \left(\frac{a\beta}{\beta^{2}} + \alpha\beta^{q} \sum_{\mu=1} dx_{\mu}^{2} \right)$$

numerical verification





• : numerical result of $d_n(x, y)$ — : analytic geodesic distance

- If x can be easily reached from y, $\theta_n(x, y)$ is small.
- $\bullet \ \theta_n(x,y) \xrightarrow{n \to \infty} 0 \quad [\forall x,y]$
- Instead of $\theta_n(x, y)$ we can also use:

 $d_n(x,y) \equiv \sqrt{-2 \ln F_n(x,y)}$ — we use this in the following

Universality of distance

Statement

There are many algorithms to construct \hat{P} satisfying (1-3) above. Among them, for a class of algorithms which generate local moves in configuration space, large scale behavior of $d_n(x, y)$ does not depend on the details of the algorithms.

(C) Geometrical optimization

•Knowing geometry, we can optimize parameters in numerical algorithms. e.g. the simulated tempering

numerical

verification

• We optimize the functional form of $\beta_a = \beta(a)$ in such a way that the distances between different modes are minimized.

This should correspond to the case in which the metric in the β direction becomes flat when *a* is used as a coordinate:

 $\frac{d\beta^2}{\beta^2} = \frac{\beta(a)^2}{\beta(a)^2} da^2 \equiv (\text{const.}) da^2.$ $\therefore \beta_a = \beta_0 e^{-Ra} \ (R : \text{const.}) \sim$ colored points: metastable states obtained in an optimization process