

# Geometric and spectral aspects of planar maps with high degrees

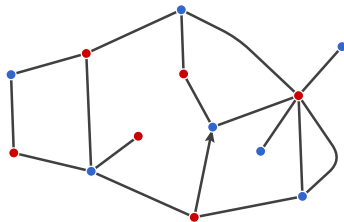
Cyril Marzouk

CNRS & Université Paris Diderot  
ERC CombiTop

Joint with  
Timothy Budd & Nicolas Curien

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- ▶ A (planar) map  $\mathfrak{m}$  is a finite connected (multi-)graph embedded in the sphere viewed up to homeomorphisms. We assume it to be rooted and bipartite.

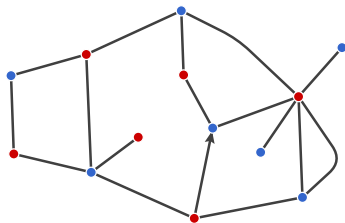


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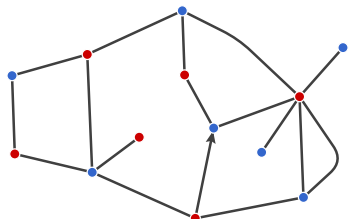


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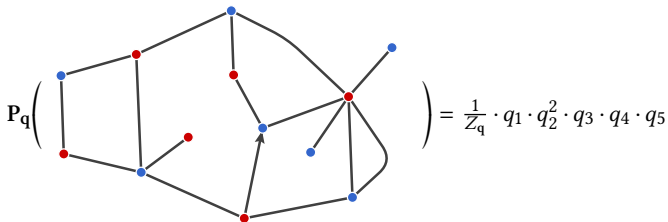
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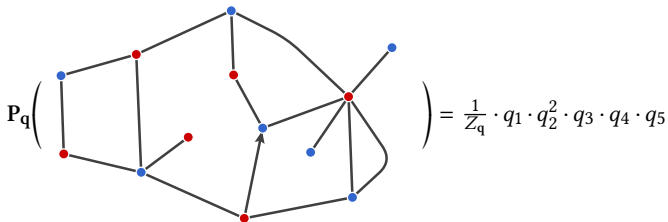
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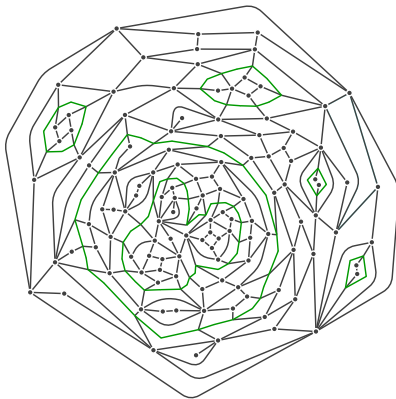
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For a typical face:  $P(\deg \geq k) \sim C \cdot k^{-(a-1/2)}$ .

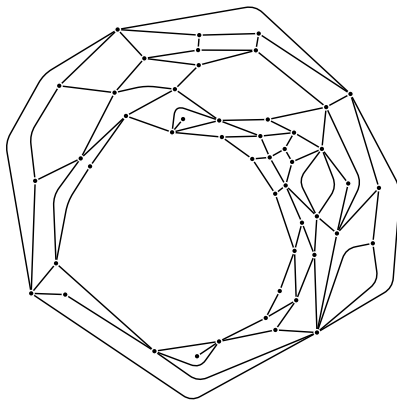




# Boltzmann maps as the gasket of a loop-decorated quadrangulation



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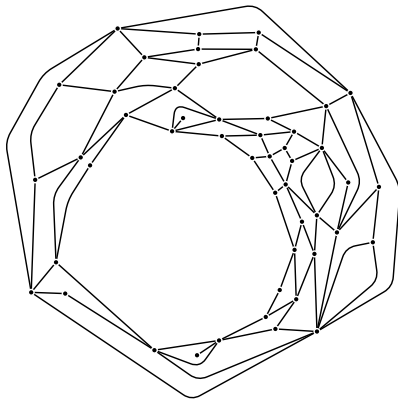


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critical  $O(n)$  loop model  $\iff$  critical Boltzmann map with

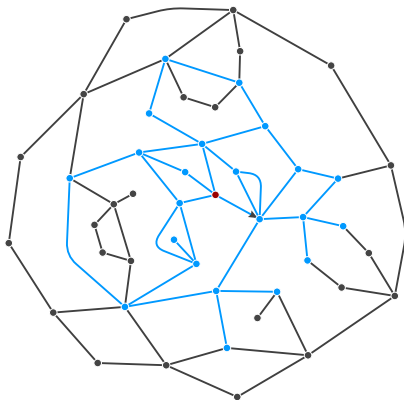
$$\begin{cases} a = 2 - \frac{1}{\pi} \arccos\left(\frac{n}{2}\right) \text{ in the } \textit{dense} \text{ phase,} \\ a = 2 + \frac{1}{\pi} \arccos\left(\frac{n}{2}\right) \text{ in the } \textit{dilute} \text{ phase.} \end{cases}$$

Supposedly related to  $\gamma$ -LQG with  $\gamma = \min(2\sqrt{a-1}, 2/\sqrt{a-1}) \in (\sqrt{2}, 2)$ .



## Balls in a map

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## Infinite Boltzmann planar maps

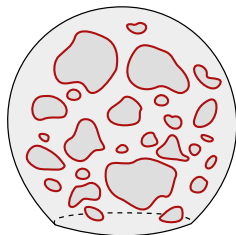
- ▶ Sample  $\mathfrak{M}_n$  from  $\mathbf{P}_q(\cdot \mid n \text{ vertices})$ , then  $\mathfrak{M}_n \rightarrow \mathfrak{M}_\infty$  in distribution in the *local limit* sense:  $\mathbf{P}(\text{Ball}_r(\mathfrak{M}_n) = \mathfrak{m}) \rightarrow \mathbf{P}(\text{Ball}_r(\mathfrak{M}_\infty) = \mathfrak{m})$  for each  $r \in \mathbf{N}$  and each map  $\mathfrak{m}$ . [Stephenson '15, also Björnberg & Stefánsson '14]

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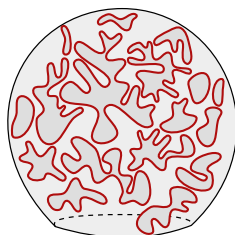
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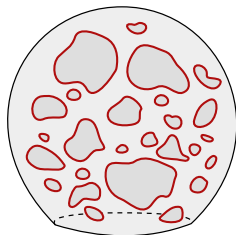
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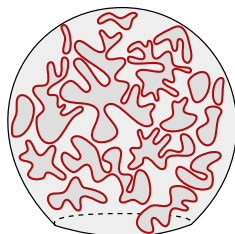
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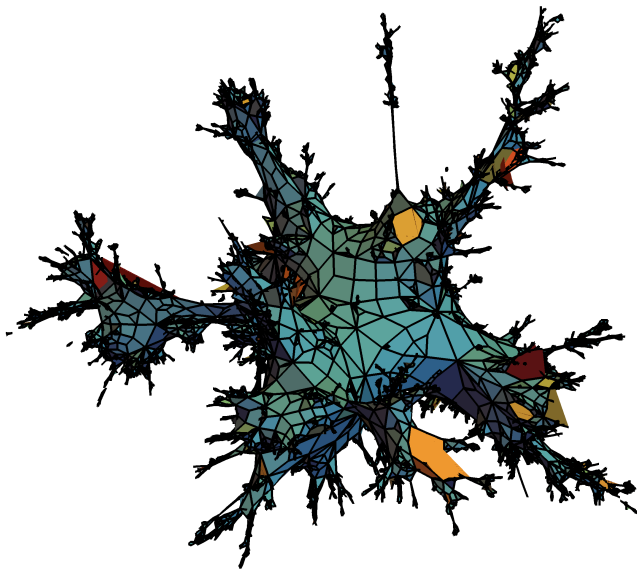


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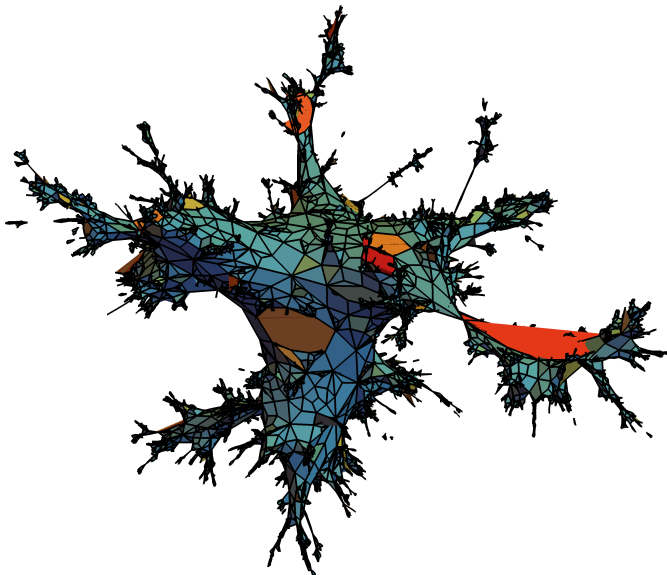
- ▶ Idea: study the *dual* map  $\mathfrak{M}_\infty^\dagger$ ! The volume  $|\text{Ball}_r(\mathfrak{M}_\infty^\dagger)|$  should increase much faster when  $a < 2$  than when  $a > 2$ .



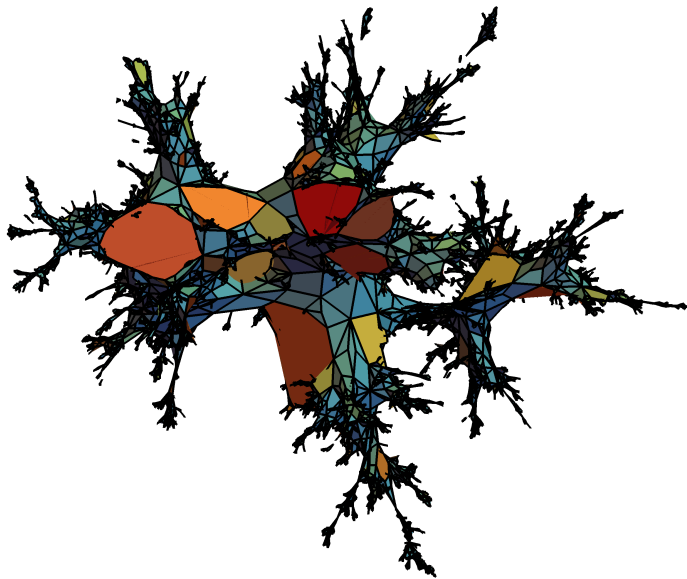
Simulation (by N. Curien):  $\mathfrak{M}_\infty$  for  $a = 2.45$



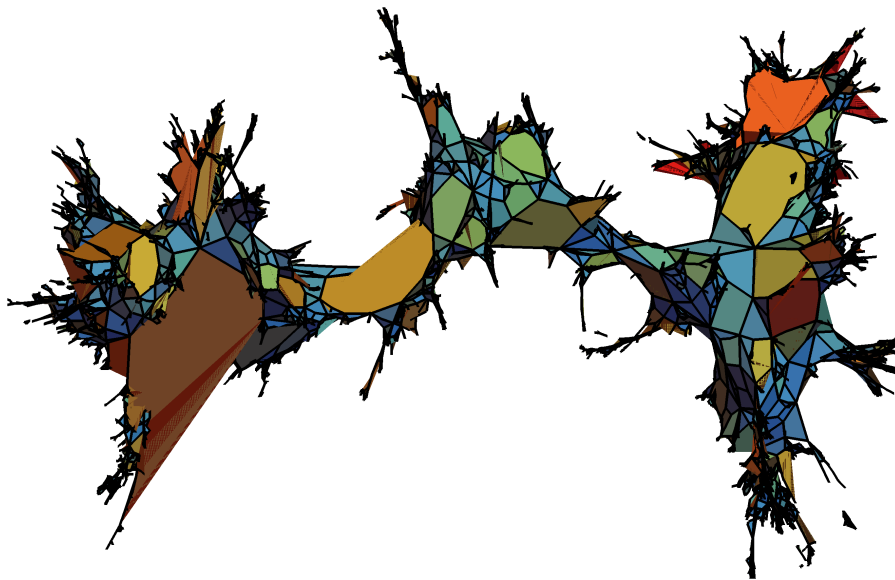
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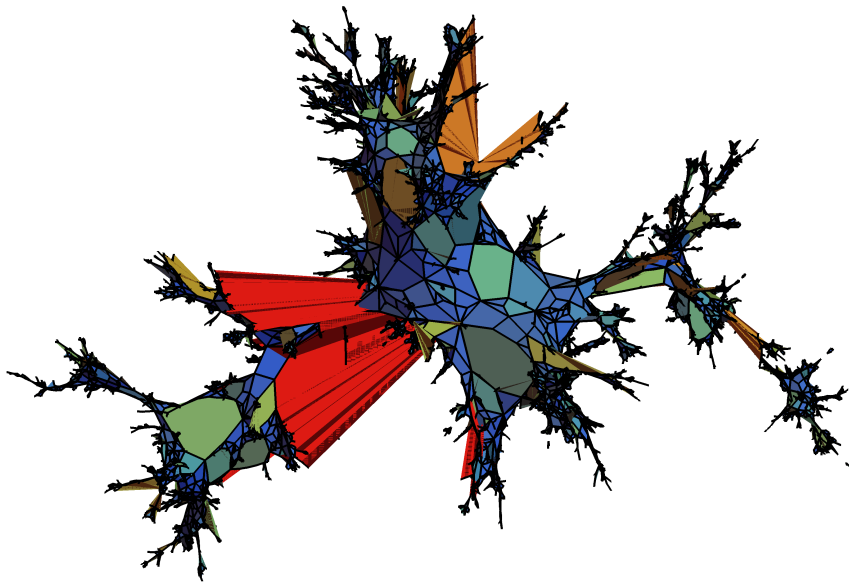
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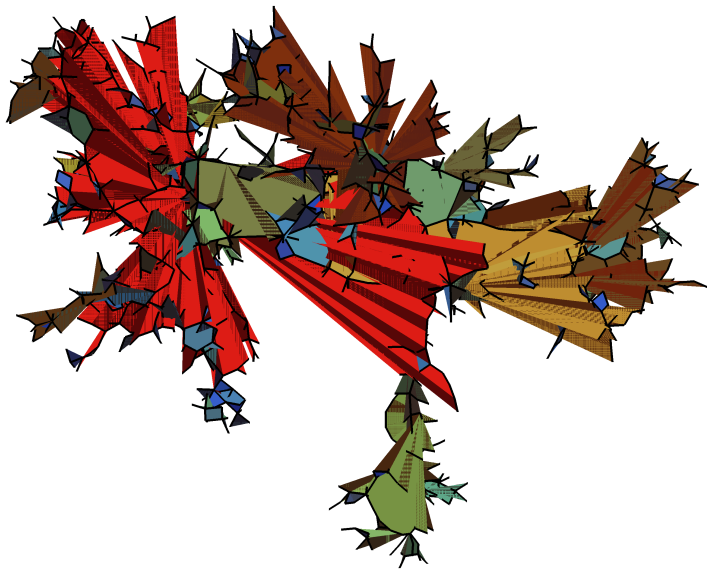
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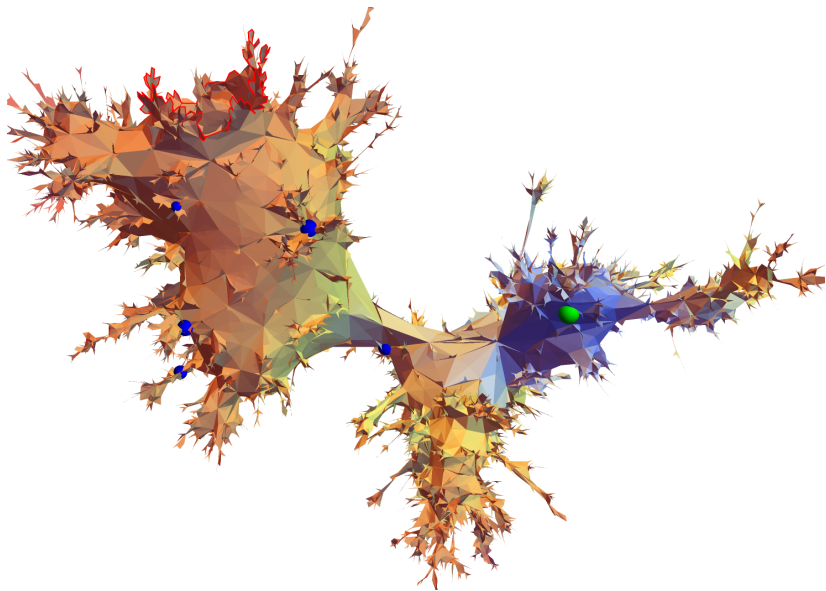
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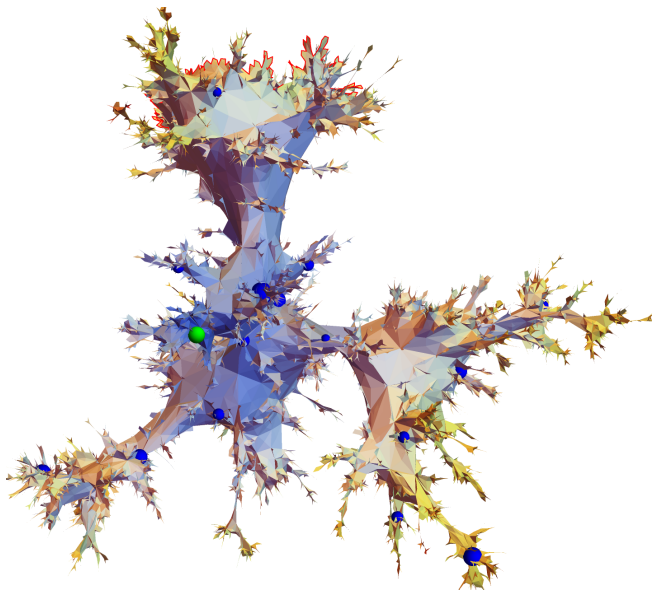


Simulation (by T. Budd):  $\mathfrak{M}_\infty^\dagger$  for  $a = 2.45$

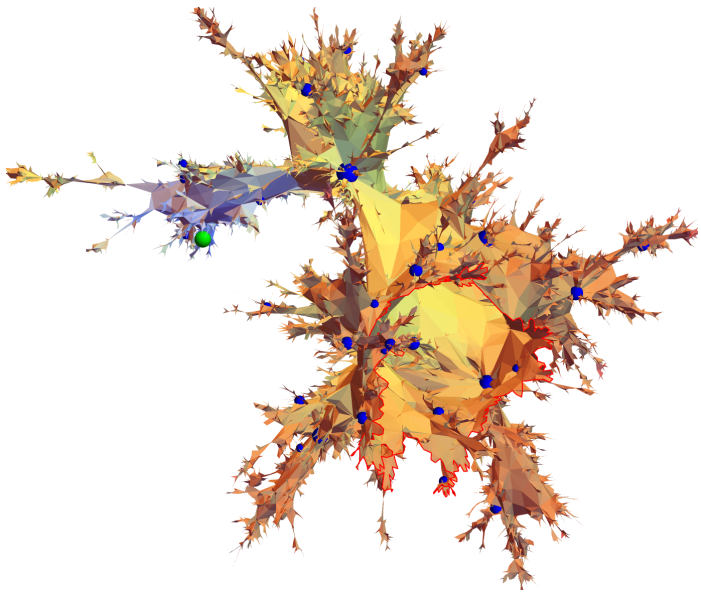




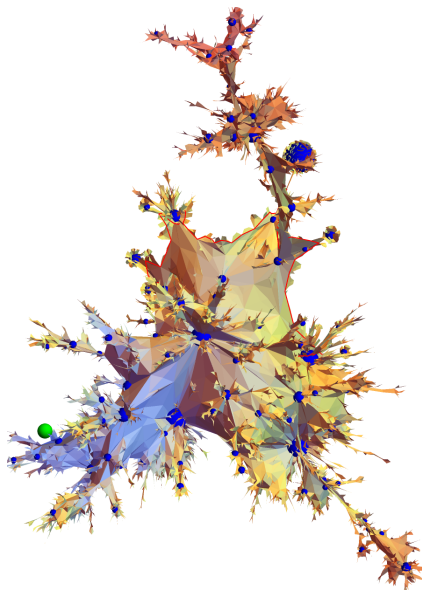
Simulation (by T. Budd):  $\mathfrak{M}_\infty^\dagger$  for  $a = 2.35$



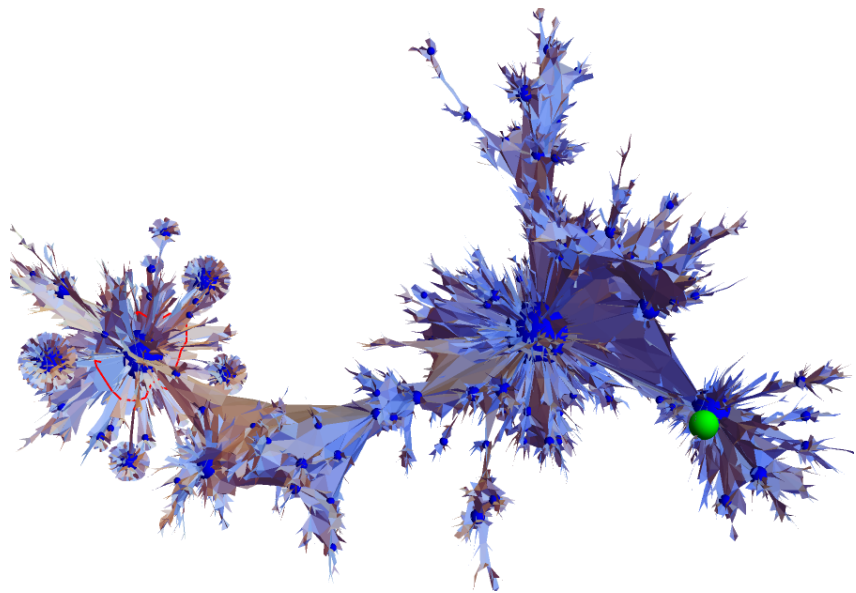
Simulation (by T. Budd):  $\mathfrak{M}_\infty^\dagger$  for  $a = 2.3$



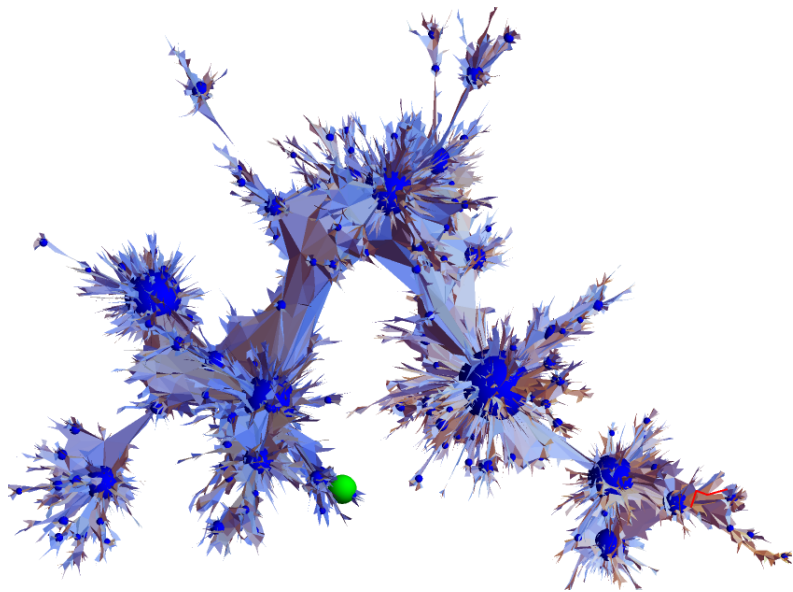
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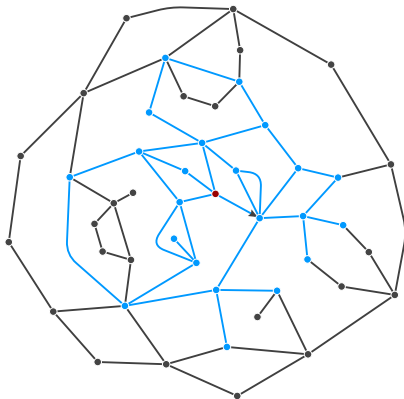


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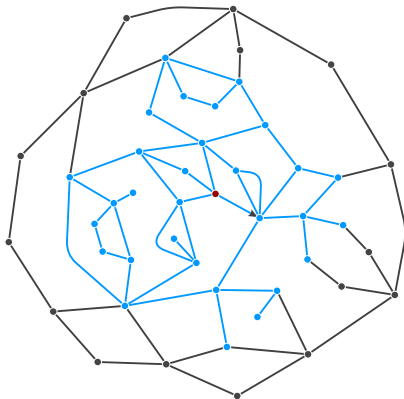
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## Volume growth

$$q_k \sim p_q \cdot c_q^{-k+1} \cdot k^{-a}$$

THEOREM (Budd & Curien '17 for  $a \neq 2$ ; Budd & Curien & ☺ '18 for  $a = 2$ ).

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In the sense:

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- ▶ If  $a < 2$ , then the shortest FPP-length of an infinite path started at the origin has finite mean!

THEOREM (Curien & ☺ in progress). For every  $a \in (\frac{3}{2}, \frac{5}{2})$ , it holds that  $|\text{Ball}_r(\mathfrak{M}_\infty)| \approx |\overline{\text{Ball}}_r(\mathfrak{M}_\infty)| \approx r^{2a-1}$ .

## Volume growth

$$q_k \sim p_q \cdot c_q^{-k+1} \cdot k^{-a}$$

THEOREM (Budd & Curien '17 for  $a \neq 2$ ; Budd & Curien & ☺ '18 for  $a = 2$ ).

	$3/2 < a < 2$	$a = 2$	$2 < a < 5/2$
$ \overline{\text{Ball}}_r(\mathfrak{M}_\infty^\dagger) $	$e^{Cr}$	$e^{3\pi(r/2)^{1/2}}$	$r^{\frac{a-1/2}{a-2}}$
$ \overline{\text{Ball}}_r^{\text{FPP}}(\mathfrak{M}_\infty^\dagger) $	$\infty$	$e^{3\pi^2 p_q r/2}$	$r^{\frac{a-1/2}{a-2}}$

In the sense:

- ▶ If  $a > 2$ , then  $(n^{-\frac{a-1/2}{a-2}} |\overline{\text{Ball}}_{[nt]}(\mathfrak{M}_\infty^\dagger)|)_{t \geq 0} \rightarrow (Z_t)_{t \geq 0}$  in distribution.
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## Random walk

$(X_n)_{n \geq 0}$  simple random walk on  $\mathfrak{M}_\infty$  and  $(X_n^\dagger)_{n \geq 0}$  simple random walk on  $\mathfrak{M}_\infty^\dagger$ .  
Recurrent or transient? diffusive?

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THEOREM (Gurel-Gurevich & Nachmias '13 for  $X$ ; Budd & Curien '17 for  $X^\dagger$ ).

	$3/2 < a < 2$	$2 \leq a < 5/2$
$(X_n)_{n \geq 0}$	recurrent	
$(X_n^\dagger)_{n \geq 0}$	transient	transient?

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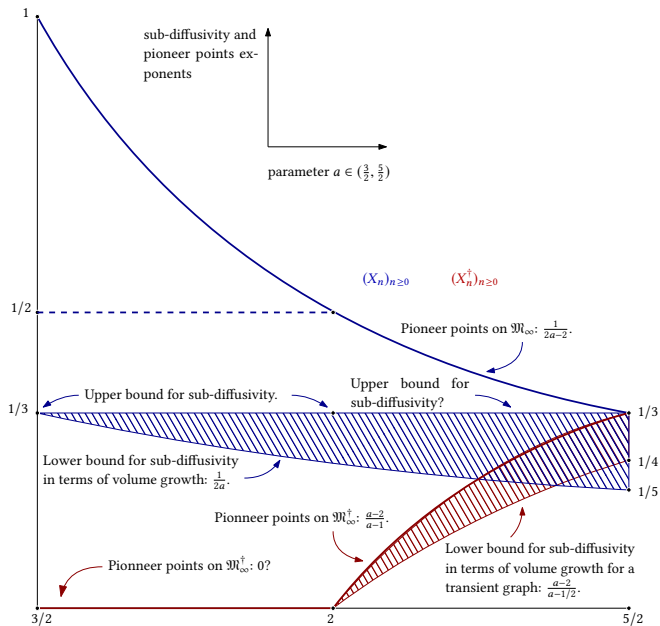
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	$3/2 < a < 2$	$2 \leq a < 5/2$
$(X_n)_{n \geq 0}$	recurrent	
$(X_n^\dagger)_{n \geq 0}$	transient	transient?

THEOREM (Curien & ☺ in progress).

	$3/2 < a < 2$	$2 \leq a < 5/2$
$d(X_n, X_0) \lesssim n^?$	1/3	$1/(2a - 2)$ ; 1/3?
$d(X_n^\dagger, X_0^\dagger) \lesssim n^?$	0?	$(a - 2)/(a - 1)$

# The big picture



## Key tool: peeling maps

# Peeling random planar maps

(Very) preliminary version

N. Curien



Gabriel Metsu (Dutch Baroque Era Painter, 1629-1667): Woman Peeling an Apple

## Key tool: peeling maps

*Y. Watabiki / Nuclear Physics B441 (1995) 119–163*

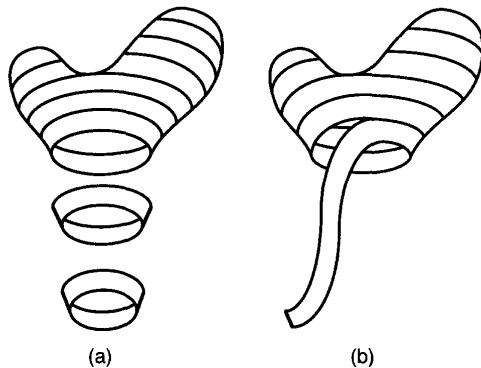
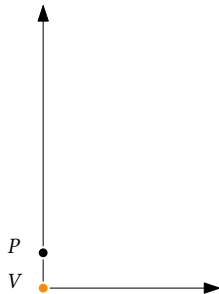
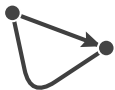
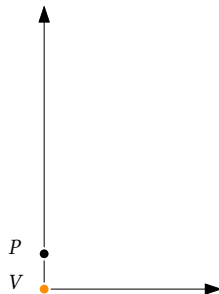
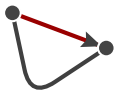


Fig. 4. Decomposition of a surface by (a) slicing and (b) peeling.

Key tool: peeling maps as Tim does

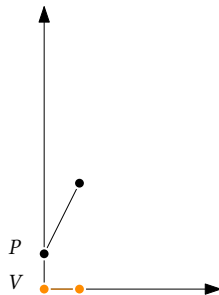
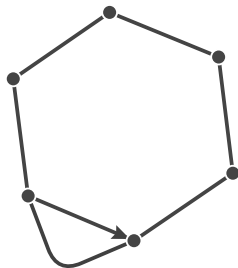


Key tool: peeling maps as Tim does

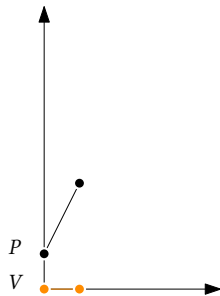
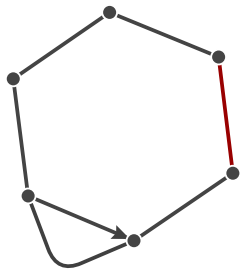




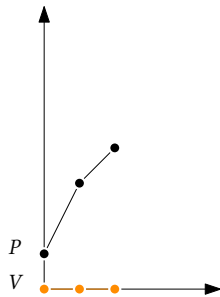
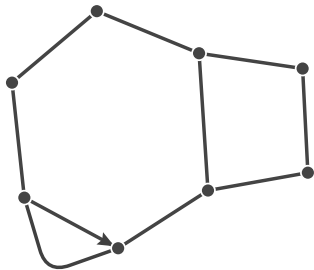
Key tool: peeling maps as Tim does



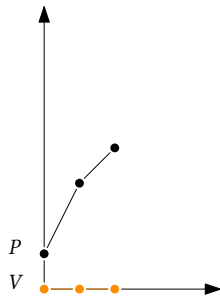
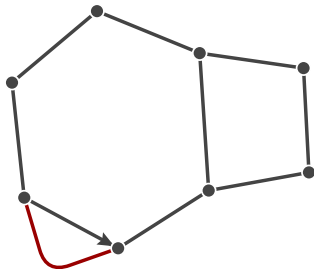
Key tool: peeling maps as Tim does



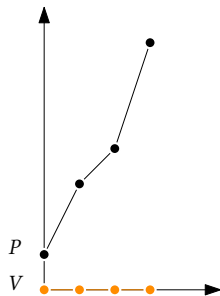
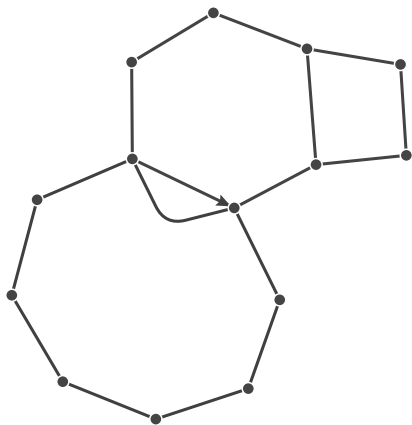
Key tool: peeling maps as Tim does



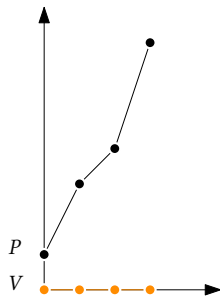
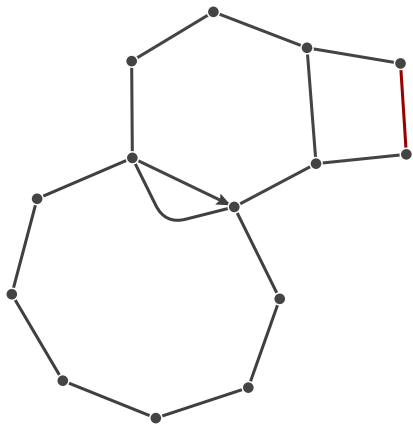
Key tool: peeling maps as Tim does



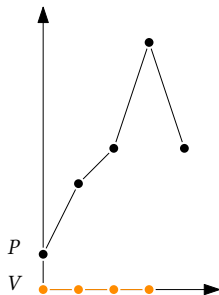
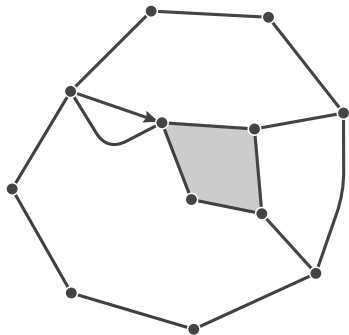
Key tool: peeling maps as Tim does



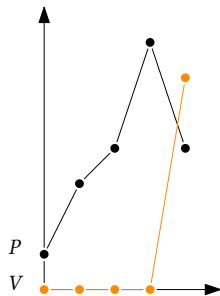
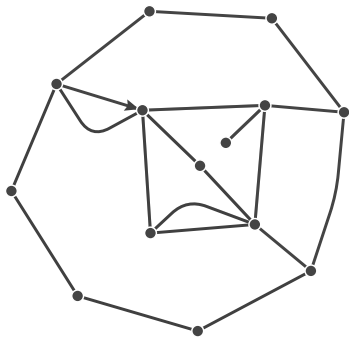
Key tool: peeling maps as Tim does



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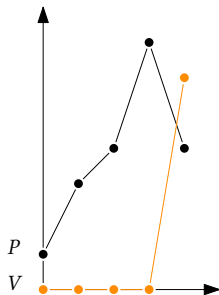
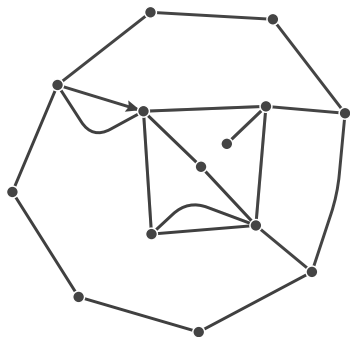


Key tool: peeling maps as Tim does





## Key tool: peeling maps as Tim does



**Budd '16:** the pair  $(P_n, V_n)_{n \geq 0}$  is a Markov chain started from  $(1, 0)$ . Precisely, from combinatorial considerations:

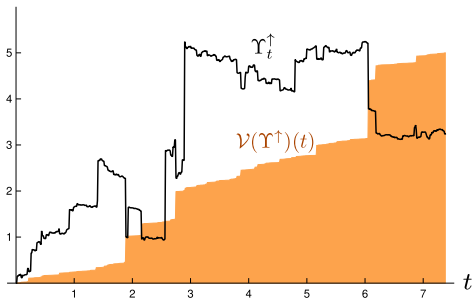
- ▶  $(P_n)_{n \geq 0}$  is some  $(a-1)$ -stable random walk *conditioned to stay positive*.
- ▶ Given  $(P_n)_{n \geq 0}$ , the increments  $(V_{n+1} - V_n)_{n \geq 0}$  are independent and if  $P_{n+1} - P_n + 1 = -\ell$ , then  $V_{n+1} - V_n$  has the law of  $|B^{(\ell)}|$ , the volume of a  $q$ -Boltzmann map conditioned to have a root-face of degree  $2\ell$ .
- ▶ **Budd & Curien '17:**  $\ell^{-(a-1/2)} |B^{(\ell)}| \rightarrow \xi$  in distribution.

## Growth in any peeling process

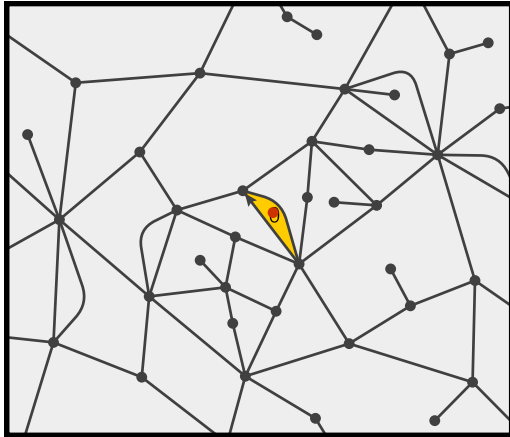
**THEOREM** (Budd & Curien '17 for  $a \neq 2$ ; Budd & Curien & ☺ '18 for  $a = 2$ ). For every  $a \in (\frac{3}{2}, \frac{5}{2})$ , it holds that

$$\left( n^{-\frac{1}{a-1}} P_{[nt]}, n^{-\frac{a-1/2}{a-1}} V_{[nt]} \right)_{t \geq 0} \xrightarrow[n \rightarrow \infty]{(d)} \left( \Upsilon^\uparrow(t), \mathcal{V}(\Upsilon^\uparrow)(t) \right)_{t \geq 0}.$$

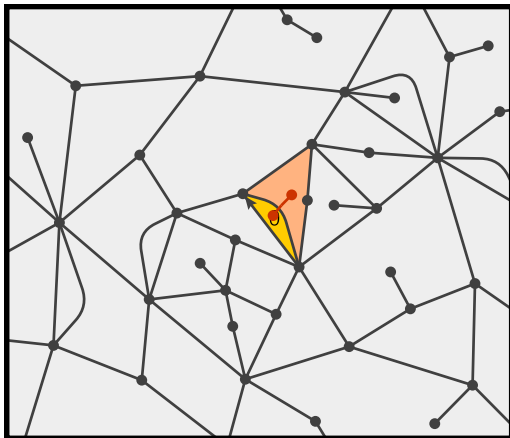
$\Upsilon^\uparrow$  is a  $(a-1)$ -stable Lévy process *conditioned to stay positive* and  $\mathcal{V}(\Upsilon^\uparrow)$  is analogous to the discrete setting.



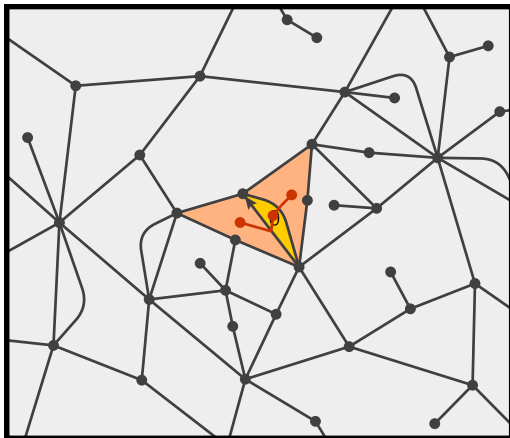
Peeling by layer



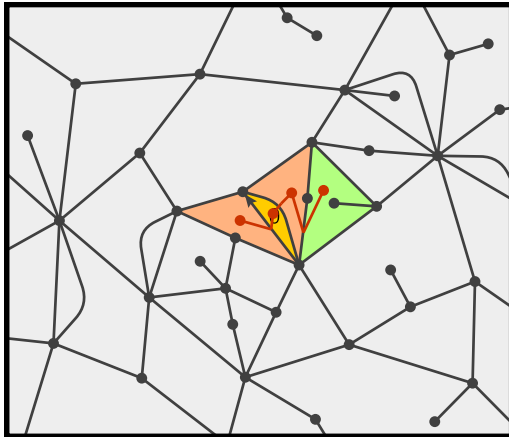
## Peeling by layer



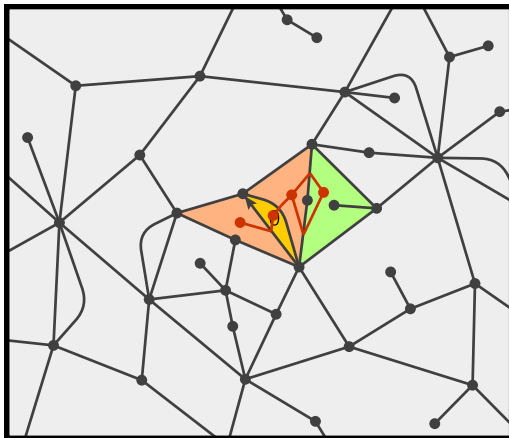
## Peeling by layer



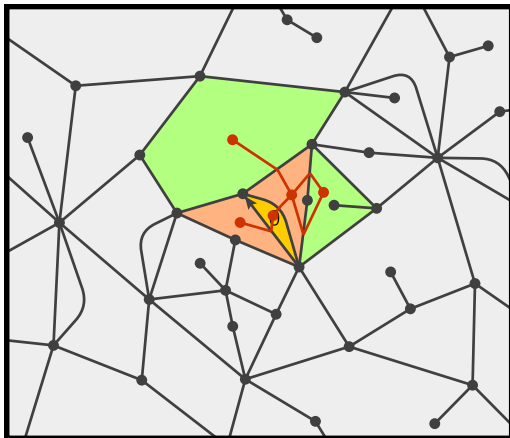
## Peeling by layer



## Peeling by layer

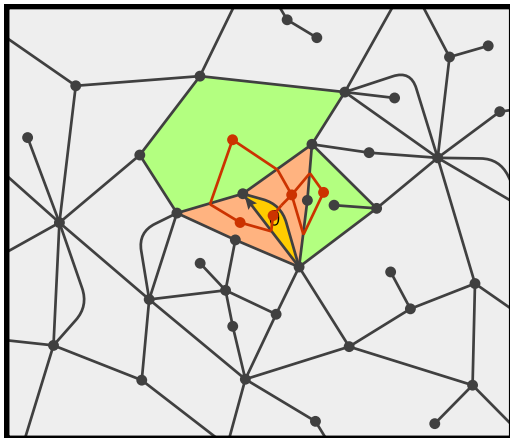


## Peeling by layer

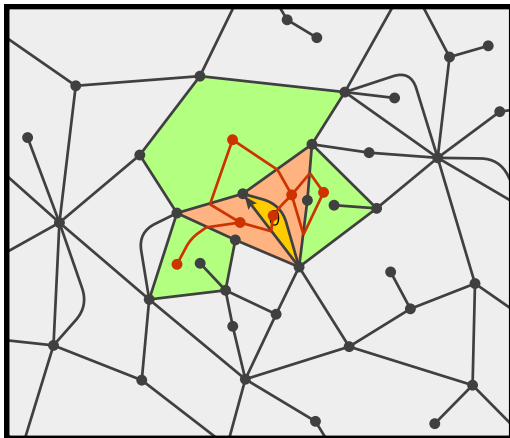




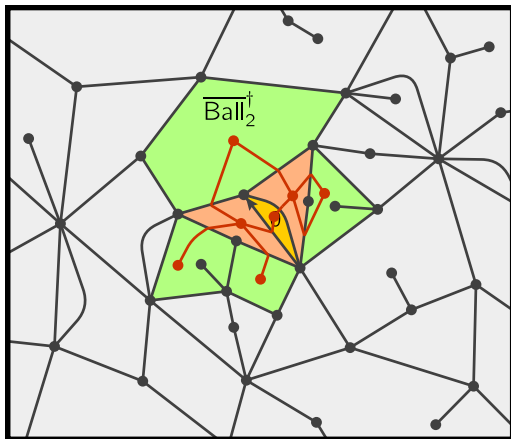
## Peeling by layer



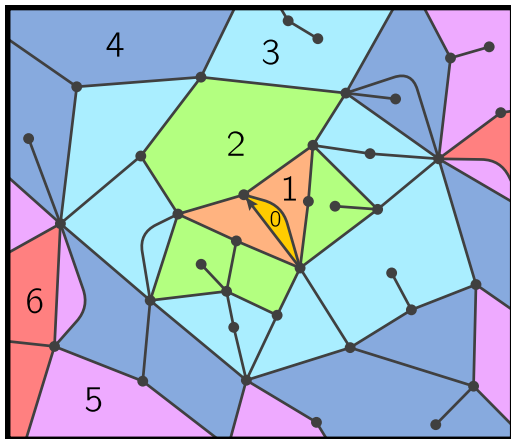
## Peeling by layer



## Peeling by layer



## Peeling by layer



## Volume growth in the dual graph distance

Can describe the dual graph distance via a peeling *by layer*.

$H_n$  = distance reached after  $n$  steps.

$\theta_r = \inf\{n : H_n \geq r\}$  = number of steps to reveal  $\overline{\text{Ball}}_r(\mathfrak{M}_\infty^\dagger)$ .

PROPOSITION.

$$\frac{H_n}{(\log n)^2} \xrightarrow[n \rightarrow \infty]{(P)} \frac{1}{2\pi^2}, \quad \text{equiv.} \quad \frac{\log \theta_r}{\sqrt{r}} \xrightarrow[n \rightarrow \infty]{(P)} \pi\sqrt{2}.$$

**Heuristic:** it takes  $\approx \frac{2\ell}{p_q \log \ell}$  steps to complete a layer of perimeter  $2\ell$ . Since  $n^{-1}P_n \rightarrow \Upsilon_1^\uparrow$ , we get

$$H_n \approx \sum_{k=0}^{n-1} \frac{p_q \log P_k}{2P_k} \approx \sum_{k=1}^n \frac{p_q \log k}{2\Upsilon_1^\uparrow k} \approx \frac{1}{\pi^2} \frac{(\log n)^2}{2}.$$

Consequently: since  $n^{-3/2}V_n \rightarrow \mathcal{V}(\Upsilon^\uparrow)(1)$ , we get

$$\frac{\log |\overline{\text{Ball}}_r(\mathfrak{M}_\infty^\dagger)|}{\sqrt{r}} = \frac{\log V_{\theta_r}}{\log \theta_r} \cdot \frac{\log \theta_r}{\sqrt{r}} \xrightarrow[r \rightarrow \infty]{(P)} \frac{3}{2} \cdot \pi\sqrt{2} = \frac{3\pi}{\sqrt{2}}.$$

## Volume growth in FPP for $a = 2$

Key observation: FPP in  $\mathfrak{M}_\infty^\dagger$  = uniform random peeling in continuous time.

$T_n$  =  $n$ -th jump-time of the process  $(|\overline{\text{Ball}}_t^{\text{FPP}}(\mathfrak{M}_\infty^\dagger)|)_t$ .

$U_r = \inf\{n : T_n \geq r\}$  = number of steps to reveal  $\overline{\text{Ball}}_r^{\text{FPP}}(\mathfrak{M}_\infty^\dagger)$ .

PROPOSITION.

$$\frac{T_n}{\log n} \xrightarrow[n \rightarrow \infty]{(P)} \frac{1}{\pi^2 p_q}, \quad \text{equiv.} \quad \frac{\log U_r}{r} \xrightarrow[n \rightarrow \infty]{(P)} \pi^2 p_q.$$

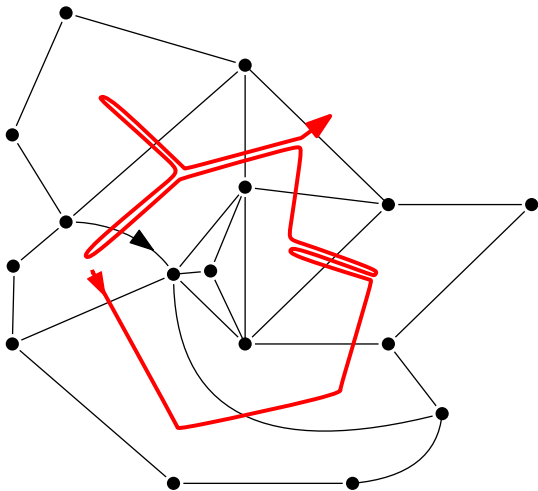
**Heuristic:**  $T_{k+1} - T_k$  is the infimum of  $2P_k$  i.i.d.  $\text{Exp}(1)$  so  $T_{k+1} - T_k \approx \frac{1}{2P_k}$ . Since,  $k^{-1}P_k \rightarrow \Upsilon_1^\uparrow$ , we get

$$T_n \approx \sum_{k < n} \frac{1}{2P_k} \approx \sum_{k < n} \frac{1}{2k\Upsilon_1^\uparrow} \approx \frac{\log n}{\pi^2 p_q}.$$

Consequently: since  $n^{-3/2}V_n \rightarrow \mathcal{V}(\Upsilon^\uparrow)(1)$ , we get

$$\frac{\log |\overline{\text{Ball}}_r^{\text{FPP}}(\mathfrak{M}_\infty^\dagger)|}{r} = \frac{\log U_r}{r} \cdot \frac{\log V_{U_r}}{\log U_r} \xrightarrow[r \rightarrow \infty]{(P)} \frac{3}{2} \cdot \pi^2 p_q = \frac{3\pi^2 p_q}{2}.$$

## Pioneer points of the RW on the dual map



## Pioneer points of the RW

$\max_{k \leq n} d(o, P_n) \lesssim n^{\frac{1}{2(a-1)}}$ ? Idea: peeling along the walk.

**For any peeling algorithm:** if  $D_n^{+/-}$  are the greatest/smallest distance from the origin to a vertex on the boundary of  $e_n$  the explored region after  $n$  peeling steps, then

$$\overline{\text{Ball}}_{D_n^-}(\mathfrak{M}_\infty) \subset e_n \subset \overline{\text{Ball}}_{D_n^+}(\mathfrak{M}_\infty).$$

Taking the volume:

$$(D_n^-)^{2a-1} \lesssim n^{\frac{a-1/2}{a-1}} \lesssim (D_n^+)^{2a-1}, \quad \text{i.e.} \quad D_n^- \lesssim n^{\frac{1}{2(a-1)}} \lesssim D_n^+.$$

Furthermore,

$$D_n^+ - D_n^- \leq \text{aper } \mathfrak{M}_\infty(\partial e_n) \lesssim |\partial e_n|^{1/2} \approx n^{\frac{1}{2(a-1)}} \quad \text{so} \quad D_n^- \approx n^{\frac{1}{2(a-1)}}.$$

**Application with the peeling along the walk:**

- ▶ Greatest distance to the origin of a pioneer point after  $n$  peeling steps  $\approx n^{\frac{1}{2(a-1)}}$ .
- ▶ Peeling only when the walk sits a pioneer point; logarithmic vertex-degrees  $\implies \approx n$  peeling steps when reaching the  $n$ -th pioneer point.



## Sub-diffusivity of the RW

$$N_R = \# \text{ steps to exit } \text{Ball}_R(\mathfrak{M}_\infty) \gtrsim R^\beta? \quad (\beta \geq 2)$$

General strategy: find a subset of edges  $\mathcal{C}_R \subset \text{Ball}_R(\mathfrak{M}_\infty) \setminus \text{Ball}_{R/2}(\mathfrak{M}_\infty)$  such that

- ▶  $|\mathcal{C}_R| \gtrsim R^\gamma$ ;
- ▶ The walk has to traverse  $\mathcal{C}_R$  to exit  $\text{Ball}_R(\mathfrak{M}_\infty)$ ;
- ▶ The walk flashed on  $\mathcal{C}_R$  is diffusive;
- ▶  $\mathcal{C}_R$  defined in an ‘ergodic’ way, so  $\mathbf{P}(k\text{-th step} \in \mathcal{C}_R) = \mathbf{P}(\text{root-edge} \in \mathcal{C}_R)$ ;
- ▶  $\mathbf{P}(\text{root-edge} \in \mathcal{C}_R) \lesssim R^{-\delta}$ ;

Then after  $N$  steps, the number of steps in  $\mathcal{C}_R$  is  $\lesssim N \cdot R^{-\delta}$ ; but when exiting  $\text{Ball}_R(\mathfrak{M}_\infty)$ , the walk has made  $\gtrsim R^{2\gamma}$  steps in  $\mathcal{C}_R$  so

$$N_R \gtrsim R^{2\gamma + \delta}.$$

In our cases,  $2\gamma + \delta = 3$ .

## Further questions

### Discrete infinite maps:

- ▶ Exact value for sub-diffusive exponent? both on  $\mathfrak{M}_\infty$  and  $\mathfrak{M}_\infty^\dagger$ .
- ▶ Transience of  $\mathfrak{M}_\infty^\dagger$  for  $a \geq 2$ ?

### Scaling limits:

- ▶  $n^{-1/(2a-1)}\mathfrak{M}_\infty \rightarrow \mathfrak{M}$  along subsequences [Le Gall & Miermont '11, ☺ '18].  
Uniqueness of the limit? 'stable map'
- ▶  $n^{-(a-1/2)/(a-2)}\mathfrak{M}_\infty^\dagger \rightarrow \mathfrak{M}^\dagger$ ? 'stable sphere'